# **Particle-base Simulations**

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# Introduction



Digitalblasphey.com







Kramer et al., Mat. Sci. Eng. A, 2008

# What? Why? How?

- Smoothed-particle hydrodynamics (SPH)
- Peridynamics
- Moving Least Squares (MLS)
- Discrete Element Method (DEM)



- Same equations apply regardless of cracks
- Cracks/damage "just happen" as a result of properties/loading
- All reduce to collection of interacting particles.
- "Meshless" = nodes, but no elements (i.e. no predefined connectivity)
- All involve intrinsic length scale (element size?)
- Involve calculation of derivatives (except for  $v = \frac{1}{4}$ )
- **Most** utilize dynamic integration of equations of motion.

# Examples: Silling's peridynamics

Inter/transgranular fracture





Cracking (plates, membranes)



Askari, et al. JOP, 2008



Macek and Silling, FE Anal. Des., 2007

Silling and Bobaru, IJNLM, 2005



# Outline

- Smoothed Particle Hydrodynamics (SPH)
- Peridynamics
- Peristatics\* with GPUs (and weird particles)
- Concluding Remarks



 Navier-Stokes: equilibrium & linear viscous fluid:

$$\rho \frac{d\underline{v}}{dt} = -\nabla p + \mu \nabla^2 \underline{v} + \underline{f}$$

• Want collection of particles to represent fluid:

$$\underline{a}_{i} = \frac{d\underline{v}_{i}}{dt} = \frac{\underline{F}_{i}}{\rho_{i}}; \quad \underline{\ddot{r}}_{i} = \frac{\underline{F}_{i}\left(\underline{r},\underline{\dot{r}}\right)}{\rho_{j}\left(\underline{r}\right)}$$

• The key to it all: (volume) smoothed (interpolated) quantities:

$$A(\underline{r}) = \int_{\Omega(h)} A(\underline{r'}) W\left(\underline{r} - \underline{r'}\right) d\underline{r}$$

$$A(\underline{r}) = \sum_{j} A_{j} V_{j} W(\underline{r} - \underline{r}_{j}) \quad V = \frac{m}{\rho} \qquad A(\underline{r}) = \sum_{j} A_{j} \frac{m_{j}}{\rho_{j}} W(\underline{r} - \underline{r}_{j})$$

• For example, density:

$$\rho_i = \rho(\underline{r}_i) = \sum_j \rho_j \left(\frac{m_j}{\rho_j}\right) W(\underline{r}_i - \underline{r}_j) = \sum_j m_j W(\underline{r}_i - \underline{r}_j)$$

100

• Kernel (interpolation) function properties:

$$\int_{\Omega(h)} W(\underline{r}) d\underline{r} = 1; \quad \lim_{h \to 0} W(\underline{r}) = \delta(\underline{r}); \quad W(\underline{r}) \ge 0$$

• *h* is critical, defines length-scale over which pressure/density are computed.

pressure 
$$W(\underline{r},h) = \frac{315}{64\pi h^9} (h^2 - |\underline{r}|^2)^3; \quad 0 \le |\underline{r}| \le h$$

viscosity 
$$W_{\mu}(\underline{r},h) = \frac{15}{2\pi h^3} \left( -\frac{|r|^3}{2h^3} + \frac{|r|^2}{h^2} + \frac{h}{2|r|} - 1 \right)$$

$$\rho \frac{d\underline{v}}{dt} = -\nabla p + \mu \nabla^2 \underline{v} + \underline{f}$$

• Need to compute gradients:

$$\nabla A(\underline{r}) = \sum_{j} A_{j} \left(\frac{m_{j}}{\rho_{j}}\right) \nabla W(\underline{r} - \underline{r}_{j})$$



$$\underline{f}_{i}^{p} = -\nabla p(\underline{r}_{i})$$
$$= -\sum_{j \neq i} p_{j} \left(\frac{m_{j}}{\rho_{j}}\right) \nabla W(\underline{r}_{i} - \underline{r}_{j})$$

Note:  $p_j$ ,  $\rho_j$  are functions of particle position!

 Gradients 'pushed' onto kernal functions, similar to FEA

• More accurate way to compute derivatives:

$$\rho \nabla A = \nabla(\rho A) - A \nabla \rho \rightarrow \nabla A = \frac{1}{\rho} \left( \nabla(\rho A) - A \nabla \rho \right)$$

• Force on particle due to pressure:

• Force on particle due to visocity:

$$\underline{f}_{i}^{\mu} = \mu \sum_{j} \left( \frac{m_{i}}{\rho_{j}} \right) \left( \underline{v}_{j} - \underline{v}_{i} \right) \nabla^{2} W(\underline{r}_{i} - \underline{r}_{j}) = \mu \sum_{j} \left( \frac{m_{i}}{\rho_{j}} \right) \left( \underline{\dot{r}}_{j} - \underline{\dot{r}}_{i} \right) \nabla^{2} W(\underline{r}_{i} - \underline{r}_{j})$$

$$\underline{a}_{i} = \frac{1}{\rho_{i}} \left( \underline{f}_{i}^{p} + \underline{f}_{i}^{\mu} \right)$$

$$\rho_{i} = f(\underline{r}_{K}) \quad \text{for all } K \text{ particles in } h$$

$$\underline{f}_{i}^{p} = f(\rho_{i}) = f(\underline{r}_{K}) \quad \text{for all } K \text{ particles in } h$$

$$\underline{f}_{i}^{\mu} = f(\underline{\dot{r}}_{K}, \underline{r}_{k}) \quad \text{for all } K \text{ particles in } h$$

- Surface tension must be included.
- Must have enough particles.



- For any time step, 'neighbors' must be identified.
- Collision detection is needed for boundaries.

## Smoothed Particle *Elastodynamics*: SPE

- Work through formalism for different constitutive law (different gradient terms)
- Time-stepping stepping slow: implicit algorithms slow.
- 'Neighbors' can be fixed for brittle materials (use cohesive laws to account for fracture)



## Examples of SPH for elastodynamics:



Cleary and Das, IUTAM Syposium, 2008: elastic-plastic projectile



Cleary and Das, IUTAM Syposium, 2008: brittle compression







Cleary, Prakash and Ha, J. Mat. Proc. Tech., 2006

# **Peridynamics (Silling)**

- Two 'flavors':
  - **Bond-based:** simply pair interactions,  $v = \frac{1}{4}$ .
  - **State-based:** peridynamic stress and deformation tensors allow arbitrary constitutive models.
- Note based on PDEs (distinction with SPH) but reduces to PDEs in limit of refinement.



Macek and Silling, FE Anal. Des., 2007x

<sup>1</sup>Note that the notation  $f = f(\xi, \eta)$  is somewhat ambigious. Indeed, for a given function u = u(x, t) the mathematical correct way to describe f is using the Nemytskii operator  $F : u \mapsto Fu$  with  $(Fu)(x, \hat{x}, t) = f(\hat{x} - x, u(\hat{x}, t) - u(x, t))$ .

## Peridynamics

• The peridynamic equation of motion:

$$\rho \underline{\ddot{u}}(\underline{x},t) = \int_{\Omega(h)} \underline{f} \left[ \underline{u}(\underline{x}',t) - \underline{u}(\underline{x},t), \underline{x}' - \underline{x} \right] dV_{\underline{x}'}$$

• The discretized form:

$$\rho \underline{\ddot{u}}_i = \sum_p \underline{f} \left( \underline{u}_p - \underline{u}_i, \underline{x}_p - \underline{x}_i \right) V_p$$

• A sensible statement:

$$\rho \underline{\ddot{u}}_i = \underline{f}_{ip} \left( \underline{dx}^{ip}, \underline{du}^{ip} \right) V_p$$



### Peridynamics

• Bond stretch between particle "i" and particle "p":

$$s_{ip} = \frac{|\underline{dx}^{ip} + \underline{du}^{ip}| - |\underline{dx}^{ip}|}{|\underline{dx}^{ip}|}$$

• The constitutive law:

$$|\underline{f}_{ip}| = k \cdot s_{ip}$$

• Equivalent continuum properties:

$$\kappa = \frac{\pi h^4 k}{18}; \quad v = \frac{1}{4} \qquad G_c = \frac{9s_o^2 \kappa \delta}{5}$$



### Peridynamics

- State-based model: compute peridynamic deformation and stress states
- Map motion onto deformation gradient
- Conventional stress-strain to get stress
- Map stress onto bond forces



Silling, McMat07 (his website)

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# **Peristatics\*** (with weird particles)

Rone Kwei Lim, William Pro, Professor Linda Petzold & Professor Marcel Utz (Southampton)



Nacre (abalone): Barthelat, et al., JMPS, 2007

Synthetic Al<sub>2</sub>O<sub>3</sub>/PMMA: Munch, et al., Science, 2008





Thermal barrier coating: Donohue, et al., Mat. Sci. Eng. A, 2013

\*(Most of) you are complicit.

Deformed composite:



- Small volumes of ceramics are strong.
- Small fractions of ductile phase limits compliance.
- Interlocking, ordered architecture transfers loads to bricks and diffuses damage.



Begley, et al., "Micromechanical models to guide the development of brick and mortar composites", JMPS, 2012

$$\overline{E}_{c} = \frac{\overline{E}_{c}}{\overline{E}_{b}} = \frac{2(\sinh[\kappa_{2}]\kappa_{1} - 2\sinh[(-1 + \overline{s})\kappa_{2}]\sinh[\overline{s}\kappa_{2}]\kappa_{2})}{2\sinh[\kappa_{2}](1 + \kappa_{1}) + (\cosh[\kappa_{2}] - \cosh[(-1 + 2\overline{s})\kappa_{2}])\kappa_{2}}$$



$$\kappa_1 = \frac{\overline{E}_m w}{\overline{E}_b t_1} \qquad \kappa_2 = \sqrt{\frac{(1 - \nu_m)\overline{E}_m w^2}{2\overline{E}_b t_2 h}}$$

 Subtle asymptotic limits: both small and large stiffnesses can be relevant



$$\kappa_1 \sim 0.1; \qquad \kappa_2 \sim 0.3$$

$$\bar{E}_c = \frac{\bar{E}_m w}{t} \left( 1 + \frac{w}{12h} \right)$$

Begley, et al., "Micromechanical models to guide the development of brick and mortar composites", JMPS, 2012



- Only brick displacements & rotation matter.
- All energy at the system is in the interface.
- Cohesive law describes energy of the interface in terms of relative motion.





- Derivatives are expensive and localization is tough to capture. (---, ---)
- Connectivity is nearest neighbor and parallel function calls are cheap. (+++,+++)
  - Monte Carlo minimization (direct search)
  - Calculate energy for new positions, accept lower energy states (and a small fraction of higher ones)





Rone Kwei Lim, William Pro, Professor Linda Petzold



Timing Results For Various Code Improvements

The computer version of Begley & Landes, 1972, ASTM STP 514

- Compute elastic energy (as a function of crack length, loading, geometry, etc.)
- Calculate the macroscopic energy release rate for the anisotropic elastic solid.
- Determine loads to initiate fracture, Infer initiation toughness.
- If toughness is actually a property, it should be independent of geometry and loading



bending displacements







Energy distribution: 90 degree case







"An Academic Life"

Screenplay by M.R. Begley & W. Pro

- Begley: "Plot the initiation toughness for all your specimens and loading conditions on a single plot."
- (Pause for figure.)
- Begley: "Is this it?!?! I thought you ran like 100 cases. Plot them all."
- William: "Why? They' re all pretty much identical."
- (Pause for dumfounded stare.)
- William: "Oh. Yeah. We should do that."









## Elastic brick and mortar particle modeling



# **Concluding Remarks**

Particle methods:

- Reduce complexity of meshing.
- Allow for discontinuities, geometry evolution.
- Can be adapted to a wide variety of constitutive models.
- Involve internal length scale: strongly impacts accuracy, speed.
- Calibration?!?!?